

Introduction

We investigated the dynamics of a qubit undergoing periodic evolution on a two-dimensional temporal space, that is part thermal and part unitary. Properties of such evolution can be studied via Non-Hermitian Floquet Hamiltonians that can be in the \mathcal{PT} -symmetric or \mathcal{PT} -broken phases. We map out the phase diagram of the system, along with the exceptional point (EP) contours through analytical and numerical methods. We analyse the system analytically for N -cycles to observe the stroboscopic behaviour through the effective Floquet Hamiltonian. Our results suggest that dynamics in unitary and thermal systems are a new avenue to realize EP degeneracies.

Model

$$H_1(t) = J(t)\sigma_x \quad H_2(t) = i\gamma(t)\sigma_z \quad (1)$$

This is our Non-Hermitian Floquet Hamiltonian where $J(t)$ and $\gamma(t)$ are arbitrary functions of time.

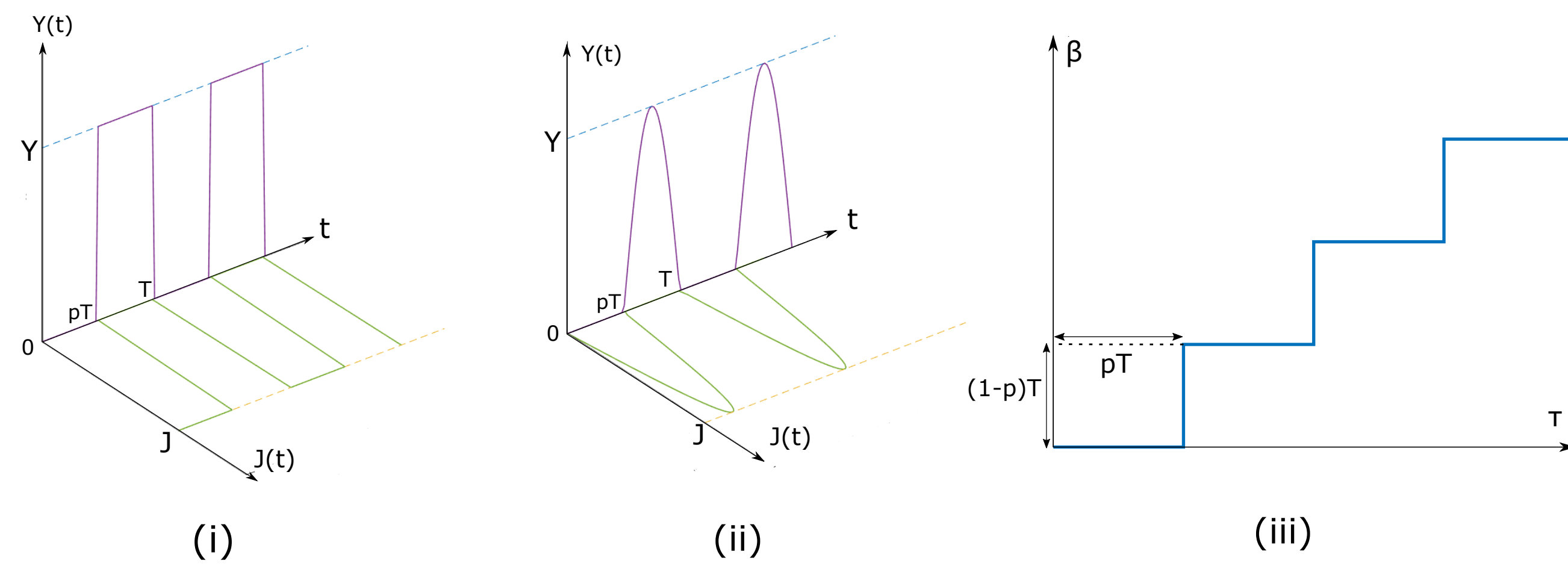


Figure 1. (i) and (ii) are representing the time-dependent Non-Hermitian Floquet model which oscillates between H_1 and H_2 with time-period T . (iii) Represents the two-dimensional temporal space where, the horizontal axis is the real time unit τ , and the vertical axis is the imaginary time unit β .

τ - Real time unit
 β - Imaginary time unit

$$\tau = \frac{pT}{2} \quad \beta = \frac{(1-p)T}{2} \quad (2)$$

U_τ - Unitary Evolution
 U_β - Thermal Evolution

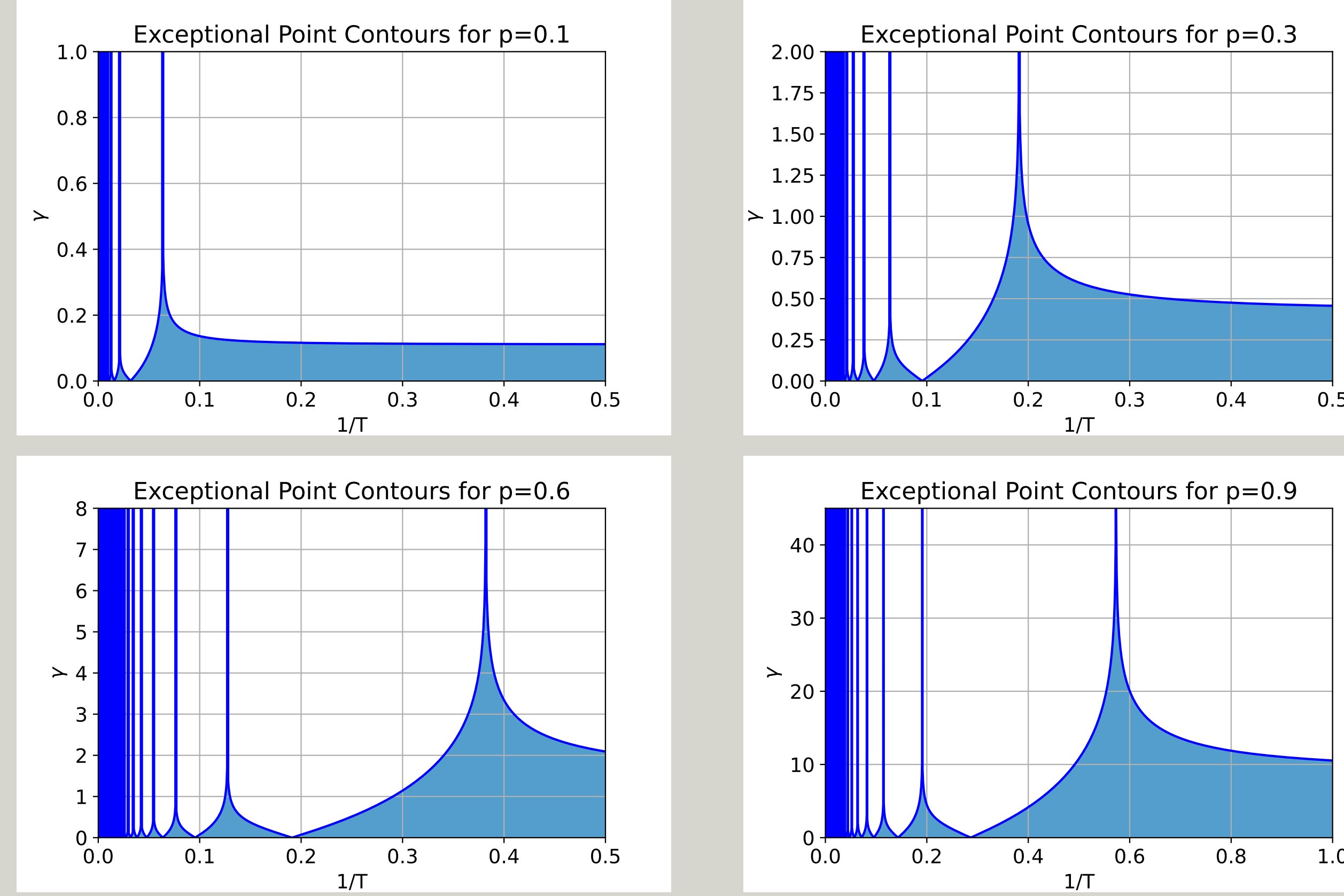
$$U_\tau = \exp\left\{-i\sigma_x \int_0^{pT} J(t)dt\right\} \quad U_\beta = \exp\left\{\sigma_z \int_{pT}^T \gamma(t)dt\right\} \quad (3)$$

Numerical Results

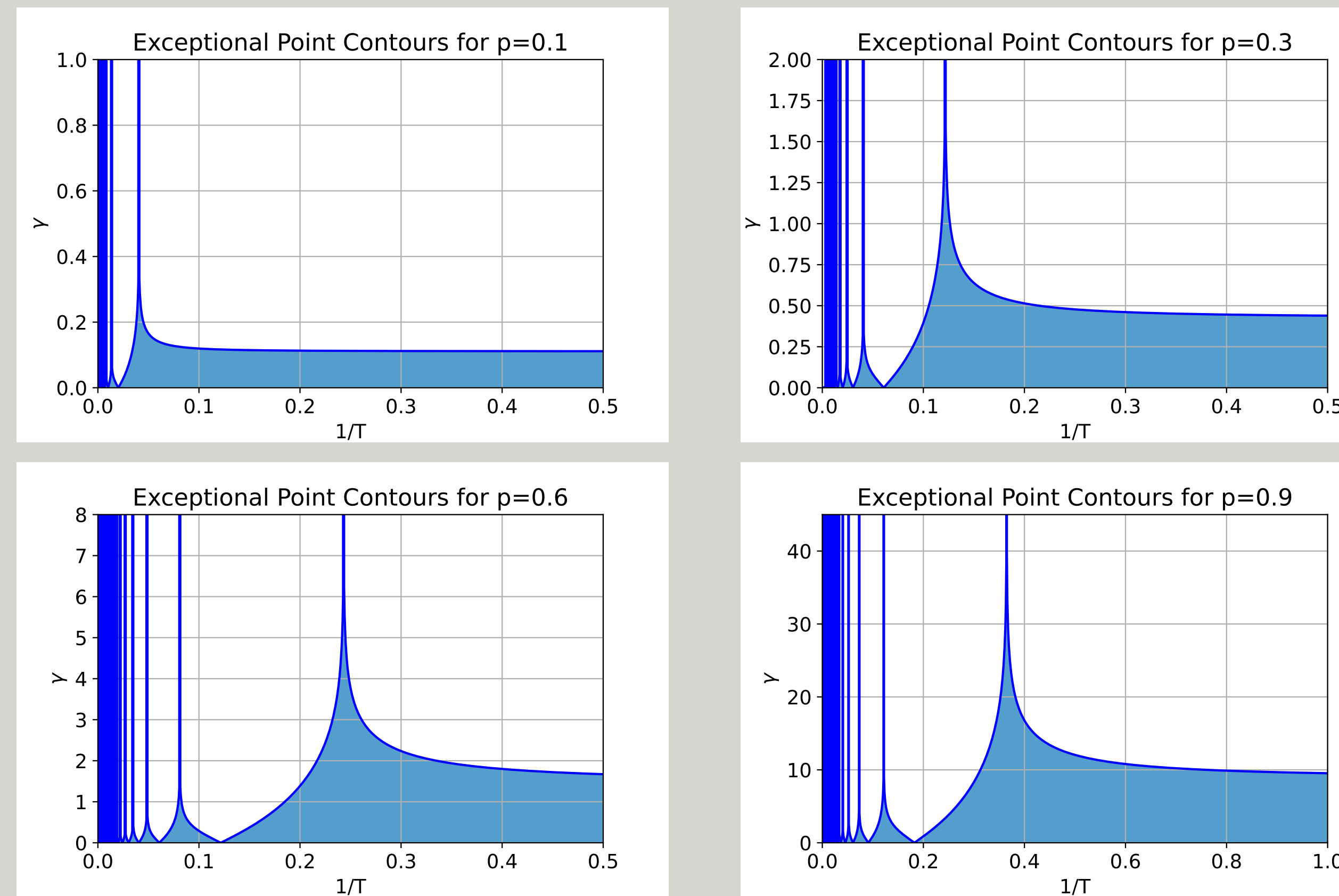
\mathcal{PT} -Symmetric region is where the eigenvalues of $U_\beta U_\tau$ are unity.

$$\cosh^2\left(\int_{pT}^T \gamma(t)dt\right) \cos^2\left(\int_0^{pT} J(t)dt\right) = 1 \quad (4)$$

Exceptional point contours (Square Wave)



Exceptional point contours (Sinusoidal Wave)



Analytical Results

- The exceptional point contours for the square wave

$$\gamma = \frac{\cosh^{-1}(|\sec(f_a p T)|)}{f_a(1-p)T} \quad (5)$$

- $T \rightarrow 0^+$

$$\gamma = \frac{p}{1-p} \quad (6)$$

- Resonance indices $\{\gamma(T) \rightarrow \infty\}$

$$T = \frac{(2k+1)\pi}{2pf_a} \quad k \in \mathbb{Z}^* \quad (7)$$

Note : f_a is defined as an areal factor, which is **identity** for Square Wave and $\frac{2}{\pi}$ for Sinusoidal Wave.

Effective Floquet Hamiltonian

The effective Floquet Hamiltonian along the EP lines

$$H_F = \frac{1}{T} \tan(2\tau J) \sigma_x + \frac{\iota}{T} \sinh(2\beta \gamma) \sin(2\tau J) \sigma_y + \frac{\iota}{T} \tanh(2\beta \gamma) \sigma_z \quad (8)$$

Evolution for N periods,

$$G(T)^N = \mathbb{1} - \iota N H_F T \equiv E_N(T) \quad (9)$$

Let $|\psi\rangle$ be the initial state, evolving it through N times,

$$\langle \psi | E_N^\dagger(T) E_N(T) | \psi \rangle = 1 + (NT)A + (N^2 T^2)B \quad (10)$$

$$A = \iota \langle \psi | H_F^\dagger - H_F | \psi \rangle$$

$$B = \langle \psi | H_F^\dagger H_F | \psi \rangle$$

Discussion

- Analysed the time-independent parameters.
- Explored the time-dependent sinusoidal behaviour.
- H_F can be used to study the stroboscopic behaviour.
- These EP degeneracies can be verified experimentally in qubit platforms.

Future Prospects

- What will happen when we couple two qubits where the first qubit is undergoing unitary evolution and the second qubit is going under thermal evolution?
For Example :

$$H = J\mathbb{1} \otimes \sigma_x + \iota \gamma \sigma_z \otimes \mathbb{1} + H_{12} \quad (11)$$

H_{12} is the interaction term.

- Will the dynamics of the system depend on the type of interaction?
- Do we expect higher order exceptional points?